

Dear Associate Editor,

Please find a revised version of our manuscript “A Double Multiple Streamtube model for Vertical Axis Wind Turbines of arbitrary rotor loading”, with changes in red, as well as detailed responses to the reviewers. We have adopted most of the reviewers’ comments in the revised manuscript. The changes include corrections in two figures, additional references and various clarifications.

Reviewer 1: Prof. Hubert Branger

We would like to thank Prof. Branger for the comments and favorable review of the manuscript. Below you will find our responses to the comments which have all been incorporated and helped improve the manuscript.

Here are some comments on A Double Multiple Streamtube model for Vertical Axis Wind Turbines of arbitrary rotor loading Az, by Anis A. Ayati, Konstantinos Steiros, Mark A. Miller, Subrahmanyam Duvvuri and Marcus Hultmark.

The double multiple stream tube model (DMST) is a worldwide analytic model used to predict the flow around vertical axis wind turbines (VAWT). DMST approach is less accurate than CFD codes, but it is far more rapid, far easier to be implemented and far more robust. The authors propose here an improvement, which allows the model to be applicable even to high loaded VAWT, i.e. with high solidity ratios and high tip speed ratios. A new momentum theory is applied to the DMST scheme. This DMST improvement is tested over direct power measurements on VAWT in the HRTF Princeton Facility. The authors modified the classical Rankine-Froude momentum theory with the new Steiros-Hultmark momentum theory, introducing the base suction effect in the wake. With this trick, the flow past the first half circle may be predicted even if the induction factor a is larger than 50% . (normally $U_{wake} = U_{input}(1-2a)$, so if a is greater than $\frac{1}{2}$, U_{wake} is null or negative...).

This new model was tested and compared to experiments performed on a very small VAWT (Radius 4.8 cm , chord length 2.1 cm, blade span: 16 cm !) but with high Reynolds numbers thanks to high density working fluids, and high solidity ratios. The new proposed DMST model provides much better power predictions than the conventional Rankine-Froude model.

It was a pleasure to read this paper. I did not see any mistake. I have just one remark about Figure 7 which seems doubtful: - Figure 7 left : the solid line is supposed to be the current model (see legend) : so why it shows negative velocity values everywhere? I thought that only the conventional model show negatives values, not the current model: I think there is a mismatch in the figure 7 legend.

Regarding figure 7a, there is indeed a mismatch in its legend: The solid line should correspond to the old model, and the points should correspond to the new model. This mistake has been corrected in the revised manuscript.

Regarding figure 7b, it is true that the two models show identical predictions, but that is expected, as figure 7b concerns the front half of the rotor. In the front half, the two momentum theories are almost identical (although one is empirical and one analytical): their inlet velocity is the free stream, while their drag predictions are almost identical, since Glauerts correction is taken into account (for induction factor greater than 0.4 for the conventional model and greater than 0.7 for the new model, see figure 2 in the manuscript). Therefore, only very slight differences are expected on the predicted quantities of the front half of the rotor, which are due to the small differences of the drag predictions of the two models (when Glauerts correction is included). The difference of predicted quantities is significant only in the rear half of the rotor, where the inlet velocity is now heavily dependent on the momentum theory choice of the front half of the rotor. This was elaborated, by adding the phrase ”The induction factor distribution is almost identical for both DMST formulations, as it concerns the front half of the rotor.” in the caption of figure 7.

However, I have to recall that static lift and drag curves have been used here, and it is known that dynamic stall plays a relevant role in the VAWT problem. Obviously there errors introduced by this fact, whatever the DMST method used. Moreover dynamic stall may be important at low TSR (i.e. in this paper) with hysteresis behavior. In the past, a lot of effort has been invested into developing modifications to the original DMST model to include those effects (Paraschivoiu

2002, Paraschivoiu and Major 1992). Most of the dynamic stall models applied to DMST consist of a series of semi-empirical procedures applied in the calculation of the lift and drag coefficients of the VAWT blade. So, I wonder what could be the performance of this new DMST model, in comparison with old fashioned DMST model but with dynamic stall corrections. I personally think that the paper can be published as it is, but a precise check of figure 7 and its legend is required, and a few sentences on DMST models with stall corrections could help the reading.

Paraschivoiu I., *Wind Turbine Design With Emphasis on Darrieus Concept*. ISBN 2-553-00931-3. Polytechnic International Press, Montreal, Canada, 2002. Ecole Polytechnique de Montreal.

Paraschivoiu I. and Major S.R., *Indicial Method calculating Dynamic Stall on a Vertical- Axis Wind Turbine.*, *Journal of Propulsion and Power*, 8(4):909-911, 1992.

It is very likely that some inaccuracies of the DMST model at low tip speed ratios (i.e. around $\lambda = 1$, see figure 9) could be a product of dynamic stall behavior, which is not taken into account in our model. Experimental investigations are ongoing, at Princeton University, in order to quantify and parameterize the dynamic stall effects on the same airfoil used for the VAWT turbine. In future work, we envision these results to be incorporated for improved DMST modeling, but it is outside of the scope of this manuscript. We have pointed out the possible effect of dynamic stall, by adding the text: “It is noteworthy to mention, however, that static data do not include the effect of dynamic stall, which is an important feature of VAWT, especially at low tip speed ratios. A better agreement of the models can be therefore expected if semiempirical corrections for dynamic stall are included in the DMST algorithm (Paraschivoiu, 2002; Major and Paraschivoiu, 1992).”, in line 233 of the manuscript, and we have also added the two references that Prof. Branger suggested.

Reviewer 2

We would like to thank the reviewer for the comments and kind words. Below you will find our responses to the comments which have all been incorporated in the revised manuscript.

The paper is well written and provides a clear description of the improved formulation of the DMST model. The reviewer has only a few suggestions:

1. In section 2.2 it is not entirely clear how the induction factor is linked to the porosity beta. Equation 4 describes CD as a function of the induction factor, while figure 2 describes CD as a function of beta. It would be helpful if the authors elaborate on this.

The method to link the open area ratio, β , and velocity through the plate is introduced in Taylor 1944 (Aero. Res. Council. R. and M. no. 2237) and expanded in Steiros and Hultmark 2018 (J. Fluid Mech. 853, R3). In summary, the method models the losses of fluid particles that pass through a pore of the plate, by assuming that all kinetic energy which is due to the acceleration of the fluid particle to enter the pore, is lost due to expansion losses. This was clarified by adding the text: “Note that in figure 2 the drag coefficient is plotted as a function of the plate porosity, β , rather than the induction factor a . These two quantities can be linked using a methodology described in the work of Steiros and Hultmark (2018), which is based on the modeling of the expansion losses of the fluid that passes through the plate” in line 112 of the manuscript.

2. In figure 3, the angle theta is indicated in the wrong corner of the velocity triangle. Consequently, $U \cos(\theta)$ and $U \sin(\theta)$ also have to be switched in the figure.

Figure 3 was corrected.

3. In figure 7, similar concerns as the other reviewer.

Regarding figure 7a, there is indeed a mismatch in its legend: The solid line should correspond to the old model, and the points should correspond to the new model. This mistake has been corrected.

Regarding figure 7b, it is true that the two models show identical predictions, but that is expected, as figure 7b concerns the front half of the rotor. In the front half, the two momentum theories are almost identical: their inlet velocity is the free stream, while their drag predictions are almost identical (although one is empirical and one analytical), since Glauert's correction is taken into account (for induction factor greater than 0.4 for the conventional model and greater

than 0.7 for the new model, see figure 2 in the manuscript). Therefore, only very slight differences are expected on the predicted quantities of the front half of the rotor, which are due to the small differences of the drag predictions of the two models (when Glauerts correction is included). The difference of predicted quantities is significant only in the rear half of the rotor, where the inlet velocity is now severely dependent on the momentum theory choice of the front half of the rotor. This was elaborated, by adding the phrase "The induction factor distribution is almost identical for both DMST formulations, as it concerns the front half of the rotor." in the caption of figure 7.

4. line 220 p11 , there is a double the the.

This mistake was corrected.

Reviewer 3

We would like to thank the reviewer for the comments and kind words. Below you will find our responses to the comments which have all been incorporated and helped improve the manuscript.

This is a manuscript on the double multiple streamtube model for horizontal-axis wind turbines. On it, a series of experiments at the high Reynolds number test facility in Princeton show that the recently model proposed by Steiros and Hutmark performs better than the classical Rankine-Froude momentum theory.

Globally, it is an interesting manuscript and the results presented are sound for the wind energy community. Also, the introduction is very well-written and presents an excellent review of state-of-the-art analytical predictive methodologies. I think the paper is suitable for publication on WES. Nevertheless, I have a few minor remarks:

- The experimental set-up section is too short, and some important information is missing. Even if the authors refer to previous works, at least the total blockage and the turbulence intensity of the incoming flow should be reported. Furthermore, as the first reviewer commented, the pipes section is small. Previous works from the group state the blockage ratio in almost 9%, that would imply the flow near the turbine is modified.

The authors appreciate the reviewer comments regarding the experimental setup section and agree that additional details would benefit the quality of the paper. To this end, section 3 in p. 9 of the manuscript was rewritten, so as to include the basic dimensions of the turbine, inlet turbulence intensities, and wind tunnel blockage due to the presence of the turbine (8.36%).

Have the authors used any blockage correction to the incoming velocity? If not, can they quantify how measurements are affected? At least the induction factor a will be modified.

The reviewer is correct that the presence of the model will affect the apparent free-stream velocity (and hence performance) of the turbine. In an effort to mitigate this, the authors elected to construct a relatively small test-model so as to achieve the lowest blockage ratio which is mechanically feasible. No further corrections have been made to the measured performance of the model primarily because the classical correction of Glauert [1], often used for wind tunnel model blockage assumes quasi-one dimensional flow through the rotor which is only valid at low induction factors $a < 0.4$, when using this type of theory for Blade Element Momentum [2]. As evident in figure 7 of this work, the typical high-solidity VAWT operates far outside the applicability of this basic correction with values exceeding $a = 0.75$. In the future, the authors would like to explore the possibility of creating a new blockage correction specifically aimed at high solidity VAWT operation.

[1] H. Glauert, Airplane Propellers, in Aerodynamic Theory (Springer, Berlin, Germany, 1935) pp. 169-360.

[2] M.O.L. Hansen, Aerodynamics of Wind Turbines, Vol. 17 (Earthscan, 2007).

- It is not clear which value of C_d is used for the conventional model on figs 6 to 9. Is it the Glauert empirical correction?

The drag coefficient used for the conventional model is given by equation (2) in the manuscript. Glauert's correction is used for induction factors greater than 0.4. This is clarified by adding the sentence: "Equations (1), (2) and (3) form the basis of the momentum theory which is incorporated

in conventional BEM models (including the conventional DMST model used in this study).” in line 102 of the manuscript.

- Although this is not important, it is not stated if the computing time of 0.7 secs on line 200 corresponds to both methods (conventional and current).

The computing time was the same in both versions of the DMST model. This was clarified by adding the sentence: “This yielded an average run time of about 0.7 seconds per λ case, for both current and conventional DMST algorithms” in line 206 of the manuscript.

- On figure 2, what does the case $\beta \sim 0$ mean?

This is the case of a solid plate, without any porous area.

- At the end of line 130, there is a typo: streamtbues.

The typo was corrected.

A Double Multiple Streamtube model for Vertical Axis Wind Turbines of arbitrary rotor loading

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Abstract. We introduce an improved formulation of the Double Multiple Streamtube (DMST) model for the prediction of the flow quantities of Vertical Axis Wind Turbines (VAWT). The improvement of the new formulation lies in that it renders the DMST valid for any induction factor, i.e. for any combination of rotor solidity and tip speed ratio. This is done by replacing the Rankine-Froude momentum theory of the DMST, which is invalid for moderate and high induction factors, with a new momentum theory recently proposed, which provides sensible results for any induction factor. The predictions of the two DMST formulations are compared with VAWT power measurements obtained at Princeton's High Reynolds number Test Facility, over a range of tip speed ratios, rotor solidities and Reynolds numbers, including those experienced by full scale turbines. The results show that the new DMST formulation demonstrates a better overall performance, compared to the conventional one, when the rotor loading is moderate or high.

10 1 Introduction

The study of vertical-axis wind turbines (VAWTs) has received renewed attention in the last decade. There were noticeable research efforts devoted to VAWTs from the mid 1970s to the mid 1980s, primarily led by Sandia National Laboratories and NASA (Strickland, 1975, 1987; Sheldahl and Klimas, 1981; Paraschivoiu et al., 1983; Paraschivoiu, 1981; Muraca et al., 1975). The following two decades saw relatively little research activity on the topic, as it was concluded that VAWTs were more prone to fatigue, and less efficient than horizontal-axis wind turbines (HAWTs). Recently, however, it was suggested that by tightly packing VAWT in a wind farm one can achieve increased power output per land area, compared to large modern HAWT farms (Dabiri, 2011). The above, coupled with the fact that VAWTs are insensitive to wind direction, have a low center of gravity, are serviceable from the ground and offer low manufacturing and maintenance costs, have created a resurgence of interest in VAWT wind farms.

20 An important prerequisite for the successful realization of wind farms is the development of engineering flow models that combine low computational cost and sufficient accuracy, so that they can be used as design and optimization tools. In the case of HAWTs, Blade Element Momentum algorithms (BEM) have been shown to fulfill these conditions, and have subse-

quently become standard aerodynamic tools of the HAWT industry. A significant amount of research has been devoted to the development of analogous models for the case of VAWTs.

25 This is not a trivial matter, however, as the aerodynamics that govern VAWTs are inherently more complex than HAWTs. The effective angle of attack experienced by a VAWT blade section is not constant, as in the case of HAWTs, but depends on the blade's instantaneous orbital position as well as on the tip speed ratio (ratio of turbine tip to free stream velocities). In addition, at relatively low tip speed ratios a blade section may experience large and rapid variations in effective angle of attack over the course of one rotation cycle. This leads to the highly unsteady and non-linear flow phenomenon known as dynamic stall (Simão Ferreira et al., 2009; Buchner et al., 2015, 2018), which causes significant hysteresis in drag and lift forces. Lastly, depending on the tip speed ratio and rotor solidity, a blade located on the downwind rotor section may interact with its own or another blade's wake generated upwind (Kozak et al., 2016; Posa and Balaras, 2018), complicating further the VAWT response.

35 Despite these inherent complexities, a number of simplified analytical predictive methodologies have been proposed over the years (e.g. vortex, cascade, fixed wake, streamtube approaches (Islam et al., 2008; Wilson and McKie, 1980)). The streamtube, and specifically its variant, the Double Multiple Streamtube (DMST) model (Paraschivoiu, 1981; Rolin and Porté-Agel, 2018) is one of the most popular approaches, due to its low computational cost, relative robustness and easiness of implementation. In a DMST model, the flow domain is discretized into a set of adjacent streamtubes, each featuring two actuators in tandem, one in the upstream rotor half-cycle and the second in the downstream half-cycle. In that way, two important aspects of the flow physics are taken into account: the constantly changing flow conditions experienced by the blades, and the difference in perceived wind between the front and rear part of the rotor.

40 Nevertheless, such treatment of the rotor fails to model other important aspects of the flow physics: DMST assumes zero expansion of the streamtubes, and it neglects the wake-blade interaction and the effect of the downstream half of the rotor on the upstream half. For these reasons, DMST algorithms are known to fail to accurately capture the local aerodynamic loads on the rotor (Wilson and McKie, 1980; Ferreira et al., 2014); still, their "global" or mean predictions are of significant accuracy, and as a result DMST remains a popular tool in VAWT design protocols.

50 Despite its usefulness, however, DMST is inapplicable to highly loaded VAWTs, i.e. characterized by high values of rotor solidity and tip speed ratio. That is because rotor loading correlates with the induction factors of the streamtubes. At an induction factor of 50% the core of the DMST model, the "classical" momentum theory of Rankine-Froude breaks down, predicting zero wake velocity and infinite wake width. For even larger induction factors the wake velocities and wake widths assume nonphysical negative values, while drag is greatly underpredicted (Hansen, 2015).

In HAWT BEM models, this inconsistency of the momentum theory is rectified by using empirical values for the drag, the so called "Glauert's correction" (Buhl and Marshall, 2005). In the case of VAWTs, however, this is not sufficient as the wake flow quantities need to be corrected as well. That is because in a DMST solution the wake velocity of the front half-rotor determines the response of the rear half-rotor. Classical momentum theory cannot accurately predict the wake flow quantities at high induction factors. As such, DMST is typically considered valid only for weakly loaded rotors where the induction factor is smaller than 50% (Ferreira et al., 2014).

In this article, we propose a resolution to this issue by substituting the Rankine-Froude momentum theory of the DMST with the momentum theory proposed by Steiros and Hultmark (2018). This momentum theory takes into account the effect of “base suction” in the wake (i.e. low wake pressure due to dissipation and wake entrainment), which is neglected in the theory of Rankine-Froude. For low induction factors, where base suction is minimal, the predictions of the two momentum theories coincide, while for large induction factors the theory of Steiros and Hultmark produces much more realistic predictions. In that way DMST becomes valid, in principle, for any rotor loading.

To quantify the accuracy of the proposed methodology, we compare predictions of a conventional DMST model equipped with both the momentum theories of Rankine-Froude and Steiros and Hultmark, and with VAWT data acquired at Princeton’s High Reynolds number Test Facility (HRTF). The data cover a range of rotor solidities, Reynolds numbers and tip speed ratios, which involve both “weakly” and “heavily” loaded rotors, at dynamically similar conditions to field VAWTs.

The structure of the article is as follows: the most relevant steps of the DMST model are outlined in section 2. The HRTF experiments are briefly described in section 3, results are discussed in section 4, and concluding remarks are given in section 5.

2 Double-Multiple Streamtube Model

In a DMST model, the rotor is divided into a front (upstream) and rear (downstream) half-cycle. The flow through a rotor of radius R is discretized into a set of adjacent streamtubes, each featuring two actuator plates in tandem, as illustrated in figure 1. The rotor circumference is divided into $2N_{st}$ arcs of equal length, $S_{st} = R\Delta\theta$, where $\Delta\theta = \pi/N_{st}$ and N_{st} is the number of streamtubes. Each streamtube is defined by an azimuth angle θ_{st} which depicts the middle point on the rotor arc bounded by the streamtube boundaries, $S_{st} = R[\theta_{st} - \frac{\Delta\theta}{2}, \theta_{st} + \frac{\Delta\theta}{2}]$. Note that the frontal area of each actuator plate, $A_{st} = dhR\Delta\theta \sin\theta_{st}$, in which dh is the length of a blade element in the spanwise direction, is not constant. Finally, an important aspect of DMST modeling is that an upstream disk is subjected to the free stream velocity, i.e. $U_{in,f} = U_\infty$, while a downstream disk is assumed to be subjected to the wake velocity produced by the front disk, i.e., $U_{in,r} = U_{f,w}$.

Using the above simplified flow description, the DMST model is able to provide predictions based on two methodologies: the momentum theory and the aerodynamic load analysis.

2.1 Classical Momentum Theory

The momentum theory aspect of conventional BEM models (including the DMST) builds upon the classical Rankine-Froude actuator disc theory (Betz, 1920; Glauert, 1930). We consider a permeable disk which acts as a drag device slowing the free-stream velocity from U_∞ far upstream, to U_a at the disc plane, and to U_w in the wake. U_a is referred to as the induced velocity and can be expressed in terms of the induction factor a defined as

$$a = 1 - \frac{U_a}{U_\infty}. \quad (1)$$

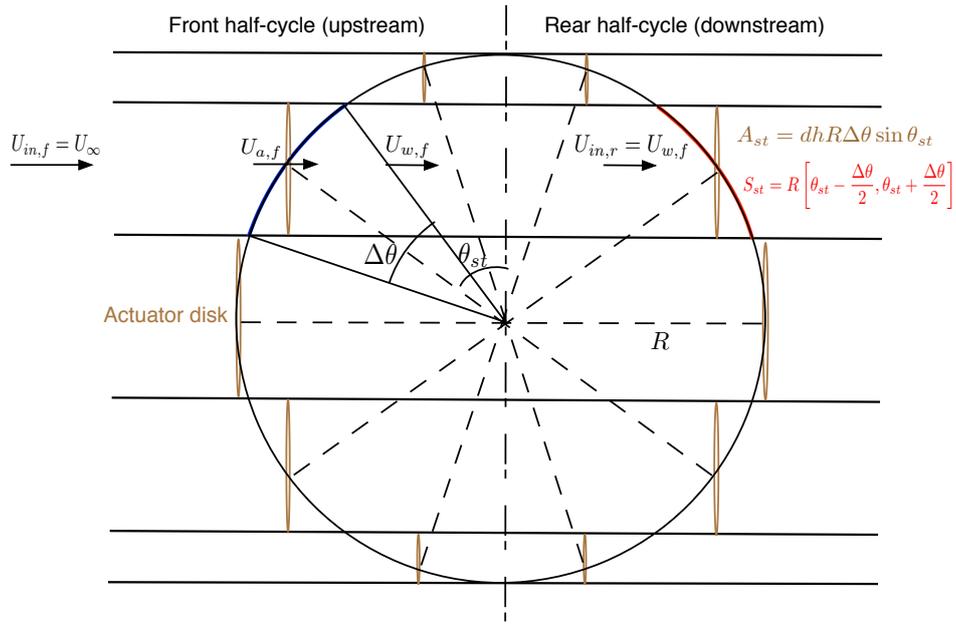


Figure 1. Schematic diagram of DMST geometrical configuration with $N_{st} = 5$.

The actuator disk theory assumes potential flow everywhere apart from the immediate vicinity of the disc, a non-rotating actuator disk and no base suction in the wake. The latter assumption implies that the wake can be treated using potential flow theory up to a point where the pressure becomes equal to the free-stream pressure, i.e. the boundary condition of the wake becomes $p_w = p_\infty$.

By applying mass and momentum balance to a control volume enclosing the actuator disk and normalizing the resulting drag with the term $\frac{1}{2}\rho AU_\infty^2$, where A and ρ are the disc area and fluid density, respectively, we obtain the well-known expression for the disk drag coefficient $C_D = 4a(1 - a)$ (Hansen, 2015). However, this expression has been shown to agree well with experimental data only for low induction factors (see figure 2). In practice, the following expression is used

$$C_D = \begin{cases} 4a(1 - a), & a \leq 0.4 \\ 0.889 - \left(\frac{0.0203 - (a - 0.143)^2}{0.6427} \right), & 0.4 < a \leq 1 \end{cases} \quad (2)$$

where the theoretical prediction is applied only for $a < 0.4$, while for larger induction factors Glauert's empirical correction (Buhl and Marshall, 2005) is used. The wake velocity $U_w(a)$ takes the form

$$U_w = U_\infty(1 - 2a). \quad (3)$$

Note that for $a > 0.5$ the momentum theory breaks down and predicts negative U_w values. The failure of the theory is even more evident if we inspect the normalized wake cross sectional area, predicted to be $A_w/A = \frac{1-a}{1-2a}$. For $a > 0.5$ this expression yields non-physical negative areas.

Equations (1), (2) and (3) form the basis of the momentum theory which is incorporated in conventional BEM models (including the conventional DMST model used in this study).

2.2 Current Momentum theory

Steiros and Hultmark (2018) extended the momentum theory of Rankine-Froude by including the effect of base suction in the wake. This theory is based on potential flow principles, where the plate is represented as a distribution of sources of equal strength. The wake velocities are rescaled, to ensure mass continuity across the plate, while the wake pressure is allowed to assume arbitrary values, so that base suction is taken into account. The various unknown quantities of the problem are then calculated using mass, momentum and energy balances.

The drag coefficient is predicted to be

$$C_D = \frac{4}{3}a \frac{(3-a)}{(1+a)}, \quad (4)$$

which, as shown figure 2, agrees well with experimental data for a larger range of plate porosities, compared to the Rankine-Froude theory. Note that in figure 2 the drag coefficient is plotted as a function of the plate porosity, β , rather than the induction factor a . These two quantities can be linked using a methodology described in the work of Steiros and Hultmark (2018), which is based on the modeling of the expansion losses of the fluid that passes through the plate.

Figure 2 shows that for low plate porosities (less than 20% of the gross plate area) the model of Steiros and Hultmark underpredicts the drag. This is because at low porosities the wake becomes unsteady and vortices are shed from the plate, a phenomenon which is not modeled by this momentum theory. However, if the wake is stabilized (e.g. with the use of a splitter plate), drag measurements collapse with the theoretical curve for all plate porosities, even up to the solid case (see figure 2).

If we express C_D as a function of the induction factor, we find that the drag predictions of Steiros and Hultmark (2018) agree well with experimental data for $a \leq 0.7$, while a correction is still needed for higher induction factors, to take into account the effect of the unsteadiness of the wake on the drag. Similarly to the classic BEM formulation, we use Glauert's empirical correction for $a > 0.7$. The resulting drag coefficient is

$$C_D = \begin{cases} \frac{4}{3}a \frac{(3-a)}{(1+a)}, & a \leq 0.7 \\ 0.889 - \left(\frac{0.0203 - (a-0.143)^2}{0.6427} \right), & 0.7 < a \leq 1 \end{cases} \quad (5)$$

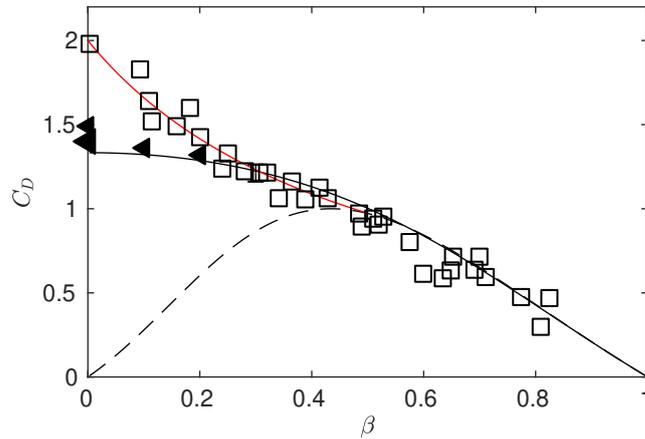


Figure 2. Porous plate drag coefficient versus plate open area ratio, $\beta = A_p/A$, where A_p is the porous plate area and A the gross plate area. Squares: measurements with no splitter plate. Triangles: measurements with splitter plate. Solid line: analytical prediction of Steiros and Hultmark (2018). Dashed line: analytical prediction of Rankine-Froude theory. Red line: Glauert's empirical correction. Adapted from Steiros and Hultmark (2018).

125 which is shown to agree with experimental measurements for all plate porosities (see figure 2). The wake velocity is predicted to be

$$U_w = \frac{1-a}{1+a}U_\infty \quad (6)$$

while the normalized wake width is predicted to be $A_w/A = 1+a$. Both wake quantities do not assume non-physical infinite or negative values at any induction factor, a fact which further demonstrates that this theory is more suitable than the Rankine-Froude theory for cases of high loading.

130 Equations 5 and 6 are used instead of equations 2 and 3 in the formulation of the novel DMST algorithm.

2.3 Aerodynamic loads analysis

The other aspect of the BEM method deals with the local aerodynamics of a blade segment (airfoil). Figure 3 provides a top-down view of a VAWT rotor and displays a velocity and force diagram associated with a blade section. The blade forces depend on the constantly changing effective angle of attack α , which is a function of the azimuth angle, θ , induction factor a and tip speed ratio $\lambda = \frac{\omega R}{U_{in}}$, where ω is the angular velocity of the turbine and $U_{in} = U_\infty$ for the front streamtubes, while $U_{in} = U_w$ for the rear **streamtubes**. From the velocity triangle, it can be seen that

$$\alpha = \arctan \left(\frac{(1-a)\sin\theta}{(1-a)\cos\theta + \lambda} \right). \quad (7)$$

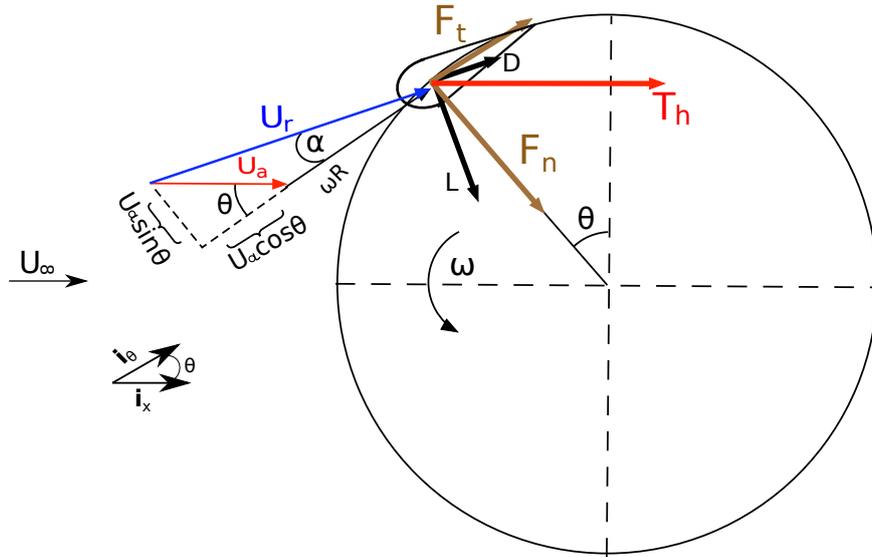


Figure 3. Velocity and force diagram on a top-down view of a VAWT rotor.

The relative velocity experienced by the blade, U_r , is the vector sum of the orbital velocity, $\omega R i_\theta$, and the induced velocity,
 140 $U_a i_x$. By virtue of equation 1 we obtain

$$U_r = U_{in} [(1-a)^2 + 2(1-a)\lambda \cos\theta + \lambda^2]^{\frac{1}{2}}. \quad (8)$$

Given the angle of attack and relative velocity, aerodynamic loads can be determined using tabulated lift and drag coefficients (C_L, C_D) and geometric considerations. In this study, static lift and drag coefficients for the airfoil profile NACA-
 145 $Re_{c_N} = [0.36, 0.70, 1.0, 2.0, 5.0] \times 10^6$ and angles of attack $\alpha \in [0^\circ, 180^\circ]$. These static Reynolds numbers are based on the free stream velocity and blade chord length. In the case of a VAWT, the static Reynolds numbers must match the “effective” blade Reynolds number based on the chord length and relative blade velocity. Subsequently, local drag and lift coefficients are found by interpolation in the $[Re, \alpha]$ -space.

It is noted that Sheldahl and Klimas (1981) did not provide measurements for the high Re quantities that we use in this
 150 study (in fact high Re airfoil data are sparse in the literature). Instead, they inferred their high- Re data using an airfoil property synthesized code which extended measurements of thinner NACA airfoil profiles, obtained at moderate Reynolds numbers. The above introduces a degree of error in the DMST predictions. Nevertheless, as shown below in the text, the predictions of the DMST model are relatively accurate for all tested Re and therefore, the inferred data of Sheldahl and Klimas (1981) can be considered reasonable estimations.

155 The drag and lift coefficients of the airfoils can be combined to yield the local tangential and normal force coefficients

$$C_t = C_D \cos \alpha - C_L \sin \alpha, \quad (9)$$

and

$$C_n = C_D \sin \alpha + C_L \cos \alpha. \quad (10)$$

By further combining C_n and C_t and multiplying with the local dynamic force $\frac{1}{2}\rho A_b U_r^2$, where A_b is the blade surface, we
 160 obtain the instantaneous thrust

$$T_h = \frac{1}{2}\rho A_b U_r^2 (C_t \cos \theta + C_n \sin \theta). \quad (11)$$

Finally, the torque τ is the product of the tangential force and radius, $\tau = F_t R$, since in our case these quantities are always orthogonal to each other, i.e.

$$\tau = \frac{1}{2}\rho A_b U_r^2 R C_t. \quad (12)$$

165 **2.4 Solving for the induction factor in a streamtube**

The DMST model calculates the induction factor a , by equating the drag of an actuator disk associated with a given streamtube to the cycle-average thrust on N_b blades that move along the rotor section S_{st} .

The cycle-average thrust coefficient corresponding to N_b blades crossing the i_{st} streamtube can be approximated as

$$C_{th}(\theta_{st}, a_{st}) = \frac{\kappa N_b \frac{1}{2\pi} \int_{\Omega_{st}} T_h(\theta, a) d\theta}{\frac{1}{2}\rho A_{st} U_\infty^2}, \quad (13)$$

170 where the integration domain is $\Omega_{st} = [\theta_{st} - \frac{\Delta\theta}{2}, \theta_{st} + \frac{\Delta\theta}{2}]$. In the limit of infinite number of streamtubes or $N_{st} \rightarrow N_\theta$, the integral in eq. (13) reduces to $\Delta\theta T_h(\theta_{st}, a_{st})$. κ is a coefficient relevant to the way blade element theory is applied in a VAWT. There is some ambiguity in the literature regarding the value of κ , which has taken different values in various streamtube algorithm implementations (e.g. $\kappa=1$ (Freris, 1990), $\kappa=2$ (Strickland, 1975) or $\kappa=4$ (Beri and Yao, 2011)). Our experimental data agree well only with the $\kappa = 4$ version, no matter the momentum theory choice; this value is therefore chosen in the DMST
 175 model and remains constant in our comparisons of current and conventional momentum theory approaches.

By equating eq. 13 to the drag of the actuator disk related to each streamtube (eq. 2 for the conventional model and eq. 5 for the new model) we obtain a nonlinear equation on a_{st} which we solve iteratively. This process is repeated twice, once for the upstream and once of the downstream rotor half-cycles.

After the induction factor a_{st} has been determined for each streamtube in both the front and rear half-cycles, the total power coefficient $C_p = C_{p,1} + C_{p,2}$, where indices 1 and 2 indicate front and rear contributions, can be computed using

$$C_p = \sum_{i=1}^2 \frac{\sum_{st=1}^{N_{st}} \frac{\kappa N_b}{2\pi} \int_{\Omega_{st}} \tau_{i,st}(\theta, a_{st}) \omega d\theta}{\frac{1}{2} \rho A_d U_\infty^3}, \quad (14)$$

where A_d is the rotor frontal area.

3 Experimental Details

In order to compare the effect of the two momentum theories in the DMST, the power of a VAWT model was tested at at Princeton's High Reynolds number Test Facility (HRTF). The HRTF is a variable-pressure, low velocity wind tunnel that can be operated at static pressures of up to $p_s = 230$ bar and free stream velocities of up to $U_\infty = 10$ ms⁻¹. This permits the testing of a large range of Reynolds numbers, while keeping the free stream velocities and Mach numbers small. The inlet turbulence intensities of the test section are 0.3% at the lowest Reynolds numbers, and 1.1% at the highest Reynolds numbers (Jiménez et al., 2010). The above facility permitted the attainment of dynamic similarity in the current experiment, by simultaneously matching the Reynolds number, tip-speed ratio and Mach number of the VAWT lab-scale models, with those encountered in full scale VAWTs.

The VAWT models (see figure 4) had a diameter $D = 96.60$ mm, a rotor span $S = 162.58$ mm, a chord $c = 21.63$ mm, leading to a tunnel blockage $SD/A_{tunnel} = 8.36\%$ (A_{tunnel} is the cross-sectional area of the tunnel's test section). A total of four lab-scale VAWT models were tested, each characterized by its number of blades ($N_b \in [2, 3, 4, 5]$). The experiments covered a range of Reynolds numbers ($5.0 \times 10^5 < Re_D < 5 \times 10^6$) and tip-speed-ratios ($0.75 < \lambda < 2.5$). Except from the number of blades, all other turbine features were identical in all four VAWT models. The airfoil profile was that of a NACA-0021. The rotor was designed to be similar to the one used in the Field Laboratory for Optimized Wind Energy (FLOWE) (Dabiri, 2011). More details related to the experimental campaign can be found in Miller et al. (2018)

4 Results and discussion

4.1 DMST convergence

In a DMST algorithm, the number of streamtubes, N_{st} , is an arbitrary parameter. To decide on that number, a convergence test was performed, based on the "normalized error"

$$\epsilon = \frac{\sum_i \sqrt{(C_{p_i} - \tilde{C}_{p_i})^2}}{\sum_i \sqrt{(C_{p_i}^2 + \tilde{C}_{p_i}^2)}}, \quad (15)$$

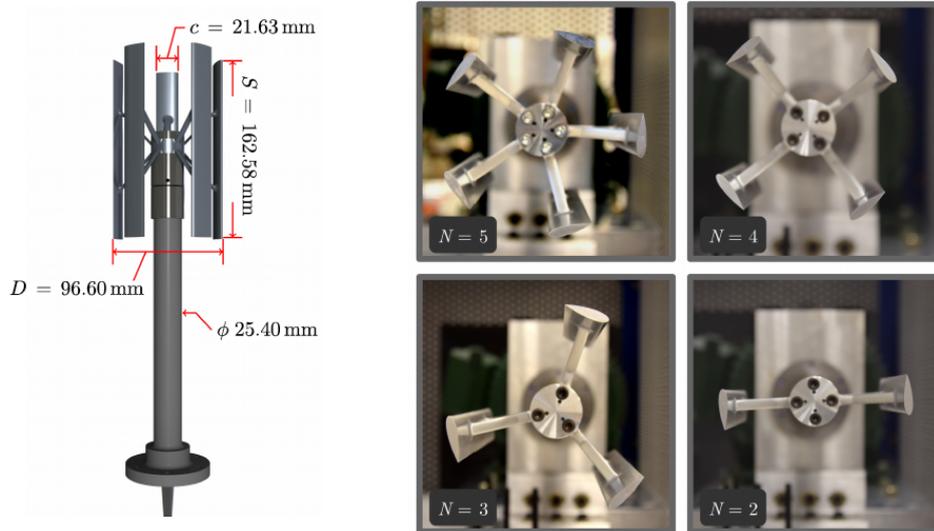


Figure 4. Left: Five-bladed VAWT model with dimensions. Right: top-down view of all four turbine configurations. The airfoil profile is that of NACA-0021 in all models.

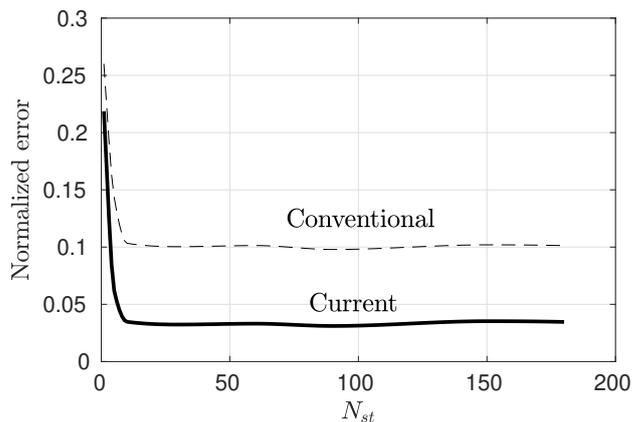


Figure 5. Normalized error as a function of number of streamtube shown for both conventional and current DMST models, for a three-bladed turbine at $Re_D = 2.85 \times 10^6$.

where Cp_i and $\tilde{C}p_i$ are the measured and predicted power coefficients for a given tip speed ratio, as indicated by the subscript i . In figure 5 we show a typical convergence plot of ϵ . The results are independent of N_{st} after approximately 15 streamtubes. We therefore used for all our tests $N_{st} = 30$. This yielded an average run time of about 0.7 seconds per λ case, **for both current and conventional DMST algorithms**, using a 3.1 GHz Intel Core i7 laptop computer running an in-house Matlab code.

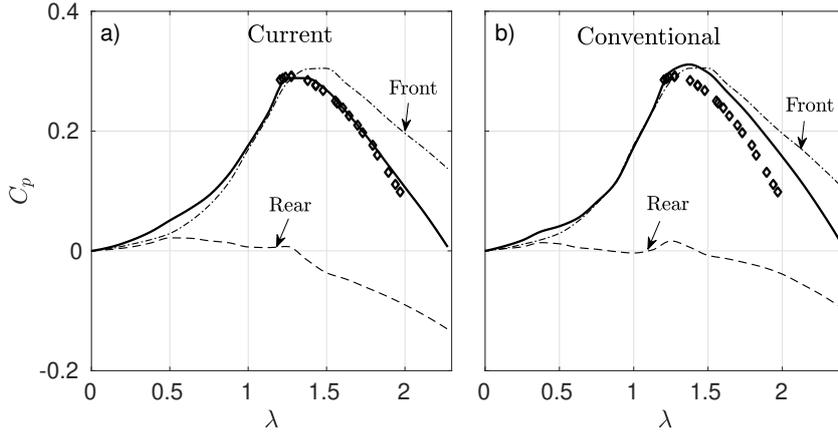


Figure 6. Comparison of current (left) and conventional (right) model predictions of power coefficients for a three-bladed VAWT. HRTF experimental data (diamonds) at $Re_D = 2.85 \times 10^6$ are plotted as a reference. Front and rear power contributions $C_{p,f}$ and $C_{p,r}$ are shown explicitly.

From figure 5 it can be qualitatively seen that the the current DMST model yields more accurate results than the conventional one. In order to assess this increase in accuracy more thoroughly, in the following sections we compare the predicted power
 210 coefficients of the two DMST versions for all four turbine configurations and across a range of operating conditions.

4.2 Experimental validation

Figure 6 shows predicted and measured power coefficients for a three-bladed VAWT. Results from the current and conventional DMST models are shown on the left and right-hand sides, respectively. The predictions include the total power coefficient C_p , and its contributions from the front ($C_{p,f}$) and rear ($C_{p,r}$) disks. The measurements correspond to tip speed ratios $1.20 < \lambda <$
 215 1.97 , free stream velocity $U_\infty = 3.1 \text{ ms}^{-1}$ and Reynolds number based on the rotor diameter $Re_D = 2.85 \times 10^6$.

As the tip speed ratio increases, the current DMST model provides power predictions which are in better agreement with the measurements, compared to the conventional one. The reason for this improvement can be seen if we compare the contributions of the front and rear disks for each model. As expected, the front power contributions are very similar, since the input velocity $U_{in,f} = U_\infty$ is the same in both models, and actuator drag is approximately captured by the Glauert correction. However, there
 220 is a noticeable difference in the rear half-cycle power predictions, due to the non-negative rear-streamtube input velocities $U_{in,r} = U_{w,f}$ of the new model.

To assess this difference in wake velocity in the above case, in figure 7 we plot the distribution of the upstream wake velocity, $U_{w,f}(\theta)$, at the highest tested tip speed ratio ($\lambda = 1.97$), that is, for the case where base suction (and therefore the difference between the two DMST implementations) is largest. We observe that, indeed, the proposed DMST model, which has the new
 225 momentum theory incorporated, predicts positive wake velocities. The conventional DMST model produces, in general, non-physical negative wake velocities. As seen from the induction factor distribution (right plot in figure 7) the negative wake

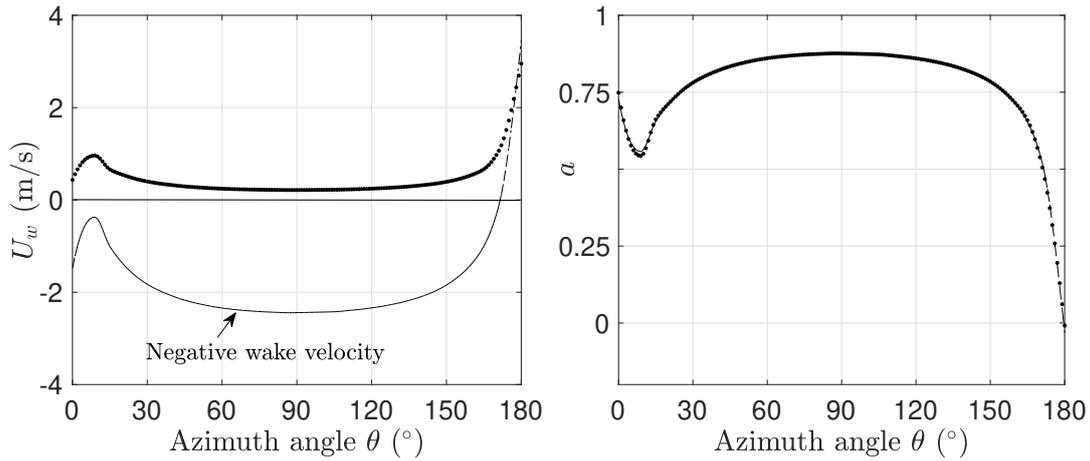


Figure 7. Front half-cycle wake velocity profile (left), U_w , and induction factor (right), a , as a function of azimuth angle θ . **Current model (points) and conventional model (solid line) at $\lambda = 1.97$. The induction factor distribution is almost identical for both DMST formulations, as it concerns the front half of the rotor.**

velocities correspond to $a > 0.5$, in accordance to the Rankine-Froude momentum theory. We remind that for such such high induction factors, DMST results based on the Rankine-Froude theory are considered invalid and were not plotted in previous studies (see for instance Ferreira et al. (2014)).

230 Figure 8 shows predicted and measured power coefficients for a three-bladed turbine at four different Reynolds numbers ($Re_D = [1.5, 2.5, 4.5, 6.0] \times 10^6$). In general, the modified DMST agrees quite well with the data and performs consistently better than the conventional DMST model at high tip speed ratios, for all Reynolds numbers. This agreement also suggests that the static airfoil data of Sheldahl and Klimas (1981), which were used in the current DMST implementation, are sufficiently accurate. **It is noteworthy to mention, however, that static data do not include the effect of dynamic stall, which is an important**

235 **feature of VAWT, especially at low tip speed ratios. A better agreement of the models can be therefore expected if semi-empirical corrections for dynamic stall are included in the DMST algorithm (Paraschivoiu, 2002; Major and Paraschivoiu, 1992).**

In figure 9 we plot the measured and predicted power coefficients for four different VAWT solidities ($N_b = 2, 3, 4$ and 5) at constant wind-tunnel conditions ($Re_D = 1.95 \times 10^6$). Again, the proposed DMST formulation improves the predictions as tip

240 speed ratio increases, for all rotor configurations.

5 Concluding remarks

A Double-Multiple Streamtube (DMST) model for vertical axis wind turbines (VAWT) is presented, where the classical Rankine-Froude momentum theory is replaced with the momentum theory of Steiros and Hultmark (2018). The classical momentum theory becomes invalid at moderate to high induction factors, and therefore limits the applicability of conventional

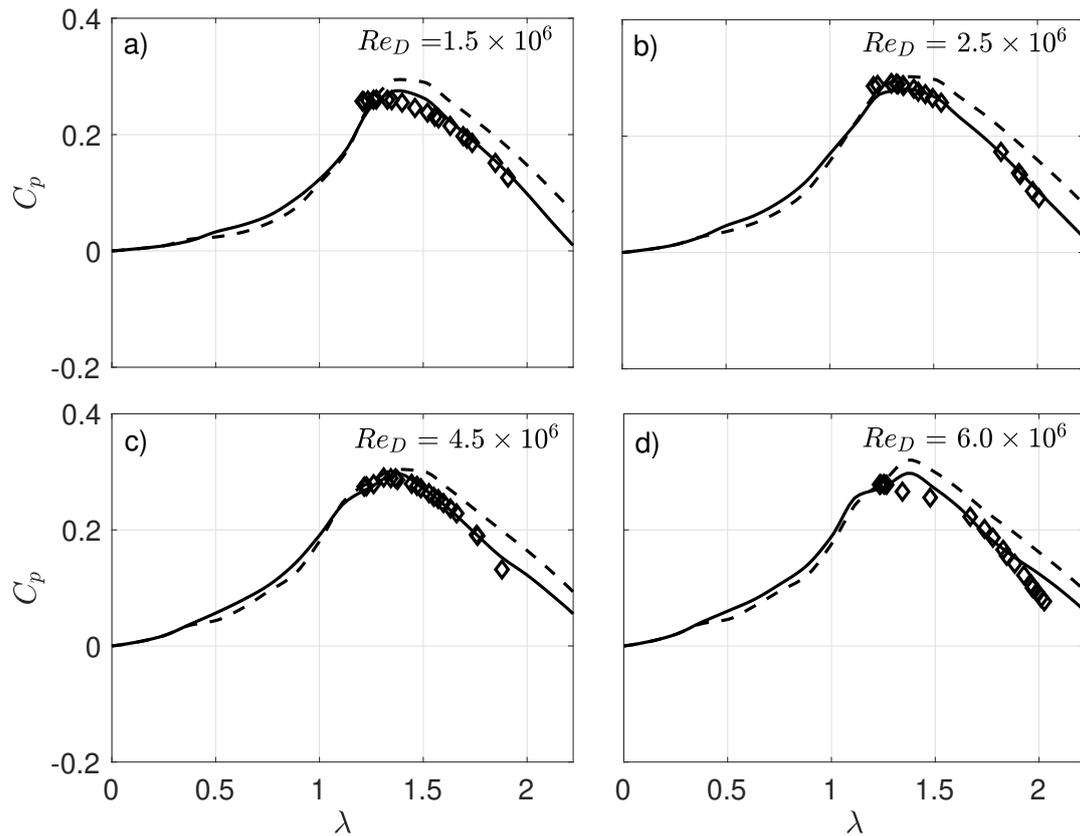


Figure 8. Measured (symbols) and predicted (current model: solid lines and conventional model: dashed lines) power coefficients, C_p , for a three-bladed VAWT ($N_b = 3$) at $Re_D = [1.5, 2.5, 4.5, 5.0] \times 10^6$.

245 DMST to small rotor solidities and tip speed ratios, that is, to small rotor loadings. In contrast, the new model introduced here is valid for any induction factor, and therefore renders the DMST applicable, in principle, to any rotor loading.

The predictions of the two DMST formulations were compared with VAWT measurements acquired at Princeton’s HRTF facility, covering a range of rotor solidities, tip speed ratios and Reynolds numbers. The data represent both lightly and heavily loaded rotors, in dynamically similar conditions to field-scale VAWTs. The results showed that the new momentum theory
 250 improves the predictions of the DMST, especially as tip speed ratio increases. It was found that this improvement is explained by a more realistic representation of the wake velocities, or equivalently input velocities to the second rear part of the rotor, from the new momentum theory.

Despite its simplicity and lack of certain flow physics, the DMST model proved reliable in its predictions of the mean power coefficient of the VAWT, for the tested range of parameters. This could be in part due to the fact that our tested tip speed ratios
 255 are rather low, while DMST inaccuracies tend to emerge at high tip speed ratios where friction and wake effects are more significant (Delafin et al., 2017).

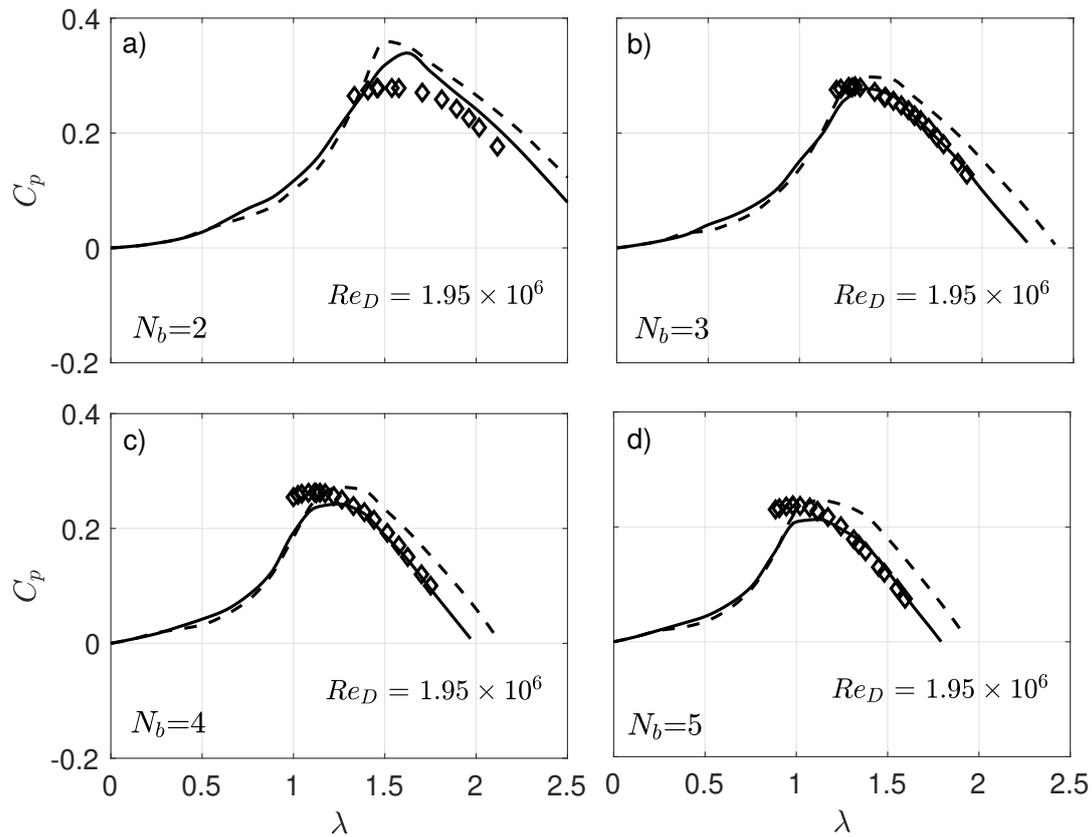


Figure 9. Measured (symbols) and predicted (current model: solid lines and conventional model: dashed lines) power coefficients, C_p , for four VAWT configurations $N_b = 2, 3, 4, 5$ and at constant $Re_D = 1.95 \times 10^6$.

Data availability. Data can be provided upon request. Please contact A. A. Ayati (ayati.anis@gmail.com).

Competing interests. The authors declare that they have no conflict of interest.

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